







Anomalous magneto-transport of Dirac-like fermions in a spin-polarized oxide two-dimensional electron system





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Secondo Congresso NQSTI, 05-07 febbraio 2025, Roma

06/02//2025









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- 1. Background: Oxide 2DES and weak (anti)localization
- 2. Anomalous magnetotransport
- 3. Nontrivial Berry curvature effect
- 4. Conclusions and Perspectives

Thank all the coauthors:

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ADVANCED MATERIALS

Dirac-Like Fermions Anomalous Magneto-Transport in a Spin-Polarized Oxide 2D Electron System

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First published: 30 October 2024 | https://doi.org/10.1002/adma.202410354

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Background: Oxide 2DES and weak (anti)localization

Letter | Published: 29 January 2004

A high-mobility electron gas at the LaAlO $_3$ /SrTiO $_3$ heterointerface

A. Ohtomo & H. Y. Hwang

Nature 427, 423–426 (2004) Cite this article LaAlO₃ (60 Å) /SrTiO₃(001) (AIO₂)а 104 (LaO)+ (AIO₂)-(□/ʊɯ) ^{XX}H (LaO)+ (TiO₂)⁰ (SrO)⁰ (TiO2)0 10-2 × 10⁻⁶ torr (SrO)⁰ 10 100

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Oxide 2DES: LaAIO₃/SrTiO₃ (LAO/STO) interfaces





- □ Mechanism for 2DES: electronic reconstruction, oxygen vacancies, cation intermixing;
- Emergent properties: Superconductivity, Ferromagnetism, Strong Rashba Spin-Orbit Coupling, tunable by gate voltage.
- Other types of 2DES:
- STO-based 2DES: cleaved STO surface, AI/STO, spin-polarized 2DES at LAO/EuTiO₃/STO interfaces...
- Non-STO system: 2DES based on KTaO₃, LAO/Ca-STO(Ferroelectric), multi-ferroic 2DES LAO/ETO/Ca-STO...









Weak localization (WL) and weak antilocalization (WAL)



Ref: Solid State Commun. 215216, 54 (2015)



Ref: Phys. Rev. Lett. 104, 126803 (2010).

- WL/WAL are the effect that conductance is suppressed/enhanced due to constructive/destructive interference between time-reversed electron self-intersecting paths.
- Dephased by magnetic field, interference is suppressed, causing positive magnetoconductance for WL while negative for WAL.



- WAL to WL with decreasing Vg in (001) LAO/STO
- Fit using Maekawa-Fukuyama (MF) formula
- Characteristic field/time/length. The extracted Rashba coefficient α_R is around 10 50 meVÅ
- SOC is strongly tunable by gate voltage









WL, WAL and beyond (competing WL and WAL)





Spin-orbit coupling Berry Phase $\gamma = \pi$

Spin-orbit coupling + Magnetic gap or surface gap opening in thin 3D-toplocial insulator Berry Phase $0 < \gamma = \pi (1 - \frac{\Delta}{2E_F}) < \pi$ Non-trivial Berry curvature and spin texture









Competing effect of WL and WAL



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Anomalous magnetotransport



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Engineering a ferromagnetic 2DES at (111)LaAIO₃/EuTiO₃/SrTiO₃ interfaces



- Epitaxial growth using Pulsed Laser Deposition (PLD) assisted with Reflected High Energy Electron Diffraction (RHEED).
- Back-gate voltage dependence of sheet resistance as a function of the temperature
- Why do we need (111) interface and ferromagnetic 2DES?

Ref: Y. Chen, at al., g, ACS Appl. Electron. Mater. 4, 3226 (2022). Y. Chen, et al., Adv. Mater. 37, 2410354 (2025).

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$(111)LaAIO_3/$ (EuTiO₃) /SrTiO₃ interfaces



- Lower symmetric: trigonal crystal field gives rise to a_{1g} and e_g^{π} derived bands.
- Spin textures: hexagonal warping and out-of-plane spin alignment.
- Berry curvature with alternatively positive and negative values at the snowflake-like Fermi contour.
- The hexagonal band-warping of (111) heterostructures unveiled large and unexpected in-plane second order bilinear magneto-resistance and anomalous Hall effect, induced by a large external in-plane magnetic field.









Magnetic Characterizations



- X-ray magnetic circular dichroism (XMCD) at Eu-M_{4,5} edge
- Anisotropic XMCD and SQUID confirm in-plane magnetization
- What is the magnetic effect on the transport behavior? Y. Chen, et al., Adv. Mater. **37**, 2410354 (2025)









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Anomalous Hall effect



- The inverse Hall coefficient decreases with increasing gate-voltage, which is opposite to the expected accumulation of electron due to the snowflake-like Fermi contour.
- Anomalous Hall effect contribution due to 2DES magnetism.

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x10









Anomalous magnetoconductance



- Anomalous shoulder-peak
- Models of Hikami-Larkin-Nagaoka, Iordanskii-Lyanda-Geller-Pikus, or Maekaewa-Fukuyama do NOT capture the MC data.
- Fitting Using formula derived in the gapped topological insulators (TIs) [Ref :*H. Z. Lu, et al, Phys. Rev. Lett. 107, 076801 (2011).*]

$$\delta\sigma(B) = \sum_{i=0,1} \frac{\alpha_i e^2}{\pi h} \left[\Psi\left(\frac{l_B^2}{l_{\phi i}^2} + \frac{1}{2}\right) - \ln(\frac{l_B^2}{l_{\phi i}^2})\right]$$

 $\alpha_0 \rightarrow WAL$ -like, $\alpha_1 \rightarrow WL$ -like Is it the ferromagnetic effect?

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Temperature dependence of magnetoconductance



- MC data as function of the temperature at Vg= -35 V, -10 V, 0 V and 5 V shown as color map and bare MC vs B data.
- Fit parameters α_i (i = 0,1) using Equation (1). The fits show that α_1 (WL-like contribution) goes to zero above 7-8 K, i.e. at a temperature similar to the FM Tc.
- The MC is associated with the gap Δ induced by ferromagnetic mangetization.

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WL and WAL at oxide interfaces



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Competing WL and WAL due to non-trivial Berry curvature



Hexagonal band warping + Rashba-like split + in-plane or canted magnetization brings about a local gap, resulting in non-trivial Berry curvature with a hot spot, analog to the MC observed in gapped TIs.

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Conclusions and Perspectives

- A method to create oxide 2DES characterized by anomalous transport properties mimicking those of systems hosting Dirac fermions, as in gapped 3D-Tls.
- Dirac-like point generated by the spin-split lowest energy bands in the simultaneous presence of Rashba-SOC, magnetic correlations, and the hexagonal symmetry of the system.
- Berry-curvature hot-spot without external planar magnetic field.

- To evaluate the Rashba Coefficient from quadratic and bilinear magnetoresistance response because we cannot use HLN, ILP, and MF formulas to fit the data.
- Second Harmonic Hall voltage detection scheme to study Berry curvature at (111) LAO/ETO/STO interfaces.
- Light control of the Non-trivial Berry Curvature : Magnetotransport under illumination with LEDs.
- spin-orbitronics and topological electronics based on the nontrivial berry curvature.











Thank you for your attention

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